

1) Notes 107(1) : Resonance in the Homopolar Generator

The homopolar generator or Faraday disk was first developed with ECE theory in pages 43 and 44. They were based on the fundamental idea of ECE theory:

$$F = A^{(0)} T \quad - (1)$$

In the case of the Faraday disk the tensor T is set up by mechanical rotation. So the basic equations of the generator are:

$$F = A^{(0)} T (\text{mechanical}) \quad - (2)$$

$$= d \wedge A + \omega \wedge A$$

and the field equations:

$$D \wedge F = R \wedge A \quad - (3)$$

$$D \wedge \tilde{F} = \tilde{R} \wedge A \quad - (4)$$

These lead to the vector equations:

$$\underline{\nabla} \cdot \underline{B} = 0 \quad - (5)$$

$$\underline{\nabla} \times \underline{E} + \frac{\partial \underline{B}}{\partial t} = \underline{0} \quad - (6)$$

$$\underline{\nabla} \cdot \underline{E} = \frac{\rho}{\epsilon_0} \quad - (7)$$

$$\underline{\nabla} \times \underline{B} - \frac{1}{c^2} \frac{\partial \underline{E}}{\partial t} = \mu_0 \underline{J} \quad - (8)$$

$$\underline{B} = \underline{\nabla} \times \underline{A} - \underline{\omega} \times \underline{A} \quad - (9)$$

$$\underline{E} = -\underline{\nabla} \phi - \frac{\partial \underline{A}}{\partial t} + \phi \underline{\omega} - \omega \underline{A} \quad - (10)$$

2) For more details of eqs (5) to (10) refer to the review paper 10 and paper 94, plus new paper by Howard E. Richard on various resonance designs in electrical engineering.

In paper 44 a complex circular basis was used to define a rotating potential set up by mechanical rotation at frequency Ω :

$$\underline{A}^{(1)} = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} - \underline{j}) e^{i\Omega t} \quad (11)$$

$$\underline{A}^{(2)} = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} + \underline{j}) e^{-i\Omega t} \quad (12)$$

This is the key concept. It is based on:

$$\underline{A} = A^{(0)} \underline{v} \quad (13)$$

so:

$$\underline{v}^{(1)} = \underline{v}^{(2)*} = \frac{1}{\sqrt{2}} (\underline{i} - \underline{j}) e^{i\Omega t} \quad (14)$$

This is rotational relativity, the mechanical spin sets up a rotational tetrad, i.e. a rotation of space-time itself.

In paper 44 a special case of eq. (10) was used:

$$\underline{E} = - \frac{\partial \underline{A}}{\partial t} - \omega \underline{A} \quad (15)$$

3) in which ϕ is assumed to be zero and where
 A is generated by the magnet of the Faraday
disk. The scalar spin current ω was assumed
to be proportional to Ω , so:

$$\underline{E}^{(2)} = \underline{E}^{(1)*} = - \left(\frac{\partial}{\partial t} + i\Omega \right) \underline{A} \quad (2)$$

$$- (16)$$

and:

$$\text{Real}(\underline{E}^{(2)}) = \text{Real}(\underline{E}^{(1)})$$

$$= \frac{2}{\sqrt{2}} A^{(0)} \Omega (i \sin \Omega t - j \cos \Omega t)$$

$$- (17)$$

This electric field strength spins around the rim
of the rotating disk. As observed experimentally
it is proportional to the product of $A^{(0)}$ and Ω .

An emf is set up between the centre and
rim of the rotating disk, and this is measured
by a voltmeter at rest w.r.t. to the disk.

Recently there have been reports of a
homopolar generator like this, exhibiting a
powerful resonance of emf as a well defined
 $\frac{1}{2}$ RPM (i.e. Ω) is approached. Re

4) a set of resonance was reported to be sudden, and resulted in an explosion of the generator. The ECE theory this is explained by spiral convection resonance.

Therefore we search for an appropriate resonance structure following the method of papers 63, 94 and the new paper a resonance design by Dr. Horst Eckhart. The simplest structure is found

from:
$$\nabla \cdot \underline{E} = \frac{f}{\epsilon_0}, \quad (18)$$

$$\underline{E} = - \frac{\partial \underline{A}}{\partial t} - \omega \underline{A}, \quad (19)$$

$$\underline{A}^{(1)} = \underline{A}^{(2)*} = \frac{A^{(0)}}{\sqrt{2}} (i - j) e^{i\Omega t} \quad (20)$$

$$\underline{E}^{(2)} = \underline{E}^{(1)*} = - \left(\frac{\partial}{\partial t} + i\Omega \right) \underline{A}^{(2)} \quad (21)$$

$$\underline{E} = \frac{2}{\sqrt{2}} A^{(0)} \Omega (i \sin \Omega t - j \cos \Omega t) \quad (22)$$

A second possibility is to use eq. (8)

with:
$$\nabla \times \underline{B} = \underline{0} \quad (23)$$

because
$$\underline{B} = B^{(0)} \underline{k} \quad (24)$$

5)

\underline{E}_0 :

$$\frac{\partial \underline{E}}{\partial t} = -c \mu_0 \underline{J} = -f_0 \underline{J} \quad (25)$$

$$\underline{E} = \frac{\partial}{\sqrt{2}} A^{(c)} \Omega (i \sin \Omega t - j \cos \Omega t) \quad (26)$$

$$\underline{E}^{(c)} = - \frac{\partial A^{(c)}}{\partial t} - i \Omega A^{(c)} \quad (27)$$

$$\frac{\partial \underline{E}^{(c)}}{\partial t} = - \frac{\partial^2 A^{(c)}}{\partial t^2} - i \Omega \frac{\partial A^{(c)}}{\partial t} = -f_0 \underline{J}$$

i.e.

$$\frac{\partial^2 A^{(c)}}{\partial t^2} + i \frac{\partial}{\partial t} (\Omega A^{(c)}) = f_0 \underline{J} \quad (28)$$

This is a resonance equation if we assume that

$$\frac{d\Omega}{dt} \neq 0 \quad (29)$$

i.e. that the angular frequency of rotation is variable. Then:

$$\frac{\partial^2 A^{(c)}}{\partial t^2} + i \Omega \frac{\partial A^{(c)}}{\partial t} + i \left(\frac{d\Omega}{dt} \right) A^{(c)} = f_0 \underline{J} \quad (30)$$

6) This is an undamped Euler-Bernoulli resonance equation if \underline{J} is periodic:

$$\underline{J} = J^{(0)} \cos(\Omega_0 t) e^{(z)} \quad (31)$$

Finally, if Ω design is such that:

$$\frac{\partial \Omega}{\partial t} \gg \Omega \quad (32)$$

then:

$$\frac{\partial^2 \underline{A}^{(z)}}{\partial t^2} + i \left(\frac{\partial \Omega}{\partial t} \right) \underline{A}^{(z)} \approx J^{(0)} \cos(\Omega_0 t) e^{(z)} \quad (33)$$

and this means that at resonance:

$$\underline{A}^{(z)} \rightarrow \infty \quad (34)$$

The explosion of the Langmuir generator may be explained in this way, i.e. using a rapidly varying Ω and a periodic current density. The latter comes from the eng set up between centre and rim of the disk, it may be an additional input.