

1) 151(2): Experimenting with New Metrics

It is simplest to begin with the Minkowski metric:

$$ds^2 = c^2 d\tau^2 = c^2 dt^2 - \underline{dr} \cdot \underline{dr} \quad - (1)$$

which is the simplest solution of the Orbital Theorem of UFT

III. The orbit for this metric is:

$$\phi(\text{Minkowski}) = \int \frac{1}{r^2} \left(\frac{1}{b^2} - \frac{1}{a^2} - \frac{1}{r^2} \right)^{-1/2} dr \quad - (2)$$

where

$$a = \frac{L}{mc}, \quad b = \frac{cL}{E} \quad - (3)$$

so

$$a = \left(\frac{mc^2}{E} \right)^{-1} b = \left(\frac{E}{mc^2} \right) b \quad - (4)$$

The Minkowski metric is:

$$\begin{aligned} \underline{dr} \cdot \underline{dr} &= c^2 (dt^2 - d\tau^2) \\ &= c^2 dt^2 \quad - (5) \end{aligned}$$

i.e.

$$\sqrt{2} dt^2 = dr^2 + r^2 d\phi^2 \quad - (6)$$

$$\sqrt{2} = \left(\frac{dr}{dt} \right)^2 + r^2 \left(\frac{d\phi}{dt} \right)^2 \quad - (7)$$

The constants of motion are

$$E = mc^2 \frac{dt}{d\tau} = \gamma mc^2 \quad - (8)$$

$$L = mr^2 \frac{d\phi}{d\tau} = \gamma mr^2 \omega \quad - (9)$$

$$\omega = d\phi / dt \quad - (10)$$

Therefore:

$$a = \frac{\gamma r^2 \omega}{c} = \gamma b \quad - (11)$$

$$b = \frac{cL}{E} = r^2 \frac{\omega}{c} \quad - (12)$$

Therefore:

$$\frac{1}{b^2} - \frac{1}{a^2} = \frac{1}{b^2} \left(1 - \frac{1}{\gamma^2} \right) \quad - (13)$$

$$= \left(\frac{v}{c} \right)^2 \frac{1}{b^2} = \frac{v^2}{r^4 \omega^2}$$

$$= \text{constant} \quad - (14)$$

So $\phi(\text{Minkowski}) = \int \frac{1}{r^2} \left(\frac{v^2}{r^4 \omega^2} - \frac{1}{r^2} \right)^{-1/2} dr$

$$\phi = \int \left(\left(\frac{v}{\omega r} \right)^2 - 1 \right)^{-1/2} dr \quad - (15)$$

where $b = \frac{r^2 \omega}{v} = \text{constant of motion} \quad - (16)$

Therefore the Minkowski metric is sufficient to give us a.s.t. Eq. (15) is equivalent to:

$$\frac{r_0}{r} < 1 \quad - (17)$$

This condition is true for light reflection by the mirror, so:

$$\Delta \phi = 2 \int_{R_0}^{\infty} \left(\left(\frac{v}{\omega r} \right)^2 - 1 \right)^{-1/2} dr \quad - (18)$$

where v is the velocity of the photon.

3) If the velocity of the photon is considered to be close to c , then:

$$\Delta\phi = 2 \int_{R_0}^{\infty} \left(\left(\frac{c}{\omega r} \right)^2 - 1 \right)^{-1/2} dr \quad - (19)$$

and the orbital angular velocity of the photon can be found from the experimental $\Delta\phi$.

In a circular orbit:

$$v = \omega r \quad - (20)$$

so the integral in eq. (18) becomes singular.