

218(7): Generalized Force Law in Polar Representation

A polar plot of the force law can be made in terms of the general $x(\theta)$ as follows:

$$F(r) = -\frac{h^2}{mr^2} \left(\frac{1}{r} + \frac{d^2}{d\theta^2} \left(\frac{1}{r} \right) \right) \quad - (1)$$

This eqn. is transformed into $F(\theta)$ using:

$$\frac{1}{r} = \frac{1}{d} (1 + \epsilon \cos(x\theta)) \quad - (2)$$

and:

$$\frac{d^2}{d\theta^2} \left(\frac{1}{r} \right) = -\frac{\epsilon}{d} \left(\left(x + \theta \frac{dx}{d\theta} \right)^2 \cos(\theta x(\theta)) + \left(2 \frac{dx}{d\theta} + \theta \frac{d^2x}{d\theta^2} \right) \sin(\theta x(\theta)) \right) \quad - (3)$$

Then for any given $x(\theta)$, F can be plotted as a function of θ is a polar plot using eq. (3).

The Special Case of Constant x .

In this case:

$$\frac{d^2}{d\theta^2} \left(\frac{1}{r} \right) = -\frac{\epsilon}{d} x^2 \cos(\theta x) \quad - (4)$$

2) and:

$$F(r) = -\frac{L^2}{mr^2} \left(\frac{1}{r} - \frac{\epsilon x^2}{d} \cos(x\theta) \right) \quad - (5)$$

where

$$d = \frac{L^2}{m^2 M G} \quad - (6)$$

So:

$$F(r) = -\frac{L^2}{mr^3} + \frac{\epsilon m M G x^2 \cos(x\theta)}{r^2} \quad - (7)$$

Here:

$$\cos(x\theta) = \frac{1}{\epsilon} \left(\frac{d}{r} - 1 \right) \quad - (8)$$

So:

$$F(r) = -\frac{L^2}{mr^3} + \frac{m M G x^2 d}{r^3} - \frac{x^2 m M G}{r^2}$$

$$F(r) = -x^2 \frac{m M G}{r^2} + (x^2 - 1) \frac{L}{mr^3} \quad - (9)$$

In polar representation:

$$F(\theta) = \frac{1}{d^2} \left(1 + \epsilon \cos(x\theta) \right)^2 \left(-x^2 m M G + \frac{(x^2 - 1) L}{m d} \left(1 + \epsilon \cos(x\theta) \right) \right) \quad - (10)$$

r. e. for constant x partial circular sections:

3)

$$F(\theta) = (1 + e \cos(x\theta))^2 \left[\frac{(x^2 - 1)L}{md^3} (1 + e \cos(x\theta)) - \frac{x^2 nMG}{d^3} \right] \quad - (11)$$

Conclusion

Any planar orbit can be described by the function $x(\theta)$ in eq. (2). The general orbit is:

$$r = f(\theta) = \frac{d}{1 + e \cos(\theta x(\theta))} \quad - (12)$$
