

261 (6): Calculation of Light Deflection due to  
gravitation from the ECE Orbital Force.

In the single polarization theory of ECE orbital  
force is:

$$\underline{F} = -\underline{\nabla} \phi - \frac{d\underline{p}}{dt} - \underline{\omega} \cdot \underline{p} + \phi \underline{\omega} \quad - (1)$$

By ECE antisymmetry:

$$-\underline{\nabla} \phi - \underline{\omega} \cdot \underline{p} = \frac{d\underline{p}}{dt} - \phi \underline{\omega} \quad - (2)$$

in which it has been assumed that:

$$\frac{d\underline{p}}{dt} = \frac{d\underline{p}}{dt} \quad - (3)$$

So:

$$\begin{aligned} \underline{F} &= 2 \left( \frac{d\underline{p}}{dt} - \phi \underline{\omega} \right) \quad - (4) \\ &= -2 (\underline{\nabla} \phi + \underline{\omega} \cdot \underline{p}) \end{aligned}$$

This result has been derived by assuming:

$$\phi_{,\mu}^a = \phi^{(0)} \eta_{\mu}^a \quad - (5)$$

2) The factor 2 in eq. (4) can be eliminated by assuming for convenience that:

$$\phi_{,\mu}^a = \frac{\phi^{(a)}}{2} v_{\mu}^a \quad - (6)$$

So:

$$\underline{F} = \frac{d\underline{p}}{dt} - \phi \underline{\omega} = -\underline{\nabla} \phi - \omega_0 \underline{p} \quad - (7)$$

As  $\omega_{\mu b}^a \rightarrow 0 \quad - (8)$

eq. (7) gives the Newtonian result:

$$\underline{F} = \frac{d\underline{p}}{dt} = -\underline{\nabla} \phi \quad - (9)$$

which is the equivalence principle:

$$\underline{F} = m \underline{g} = -mM G \frac{\underline{r}}{r^3} \quad - (10)$$

where  $\phi = -\frac{mM G}{r} \quad - (11)$

In order to calculate light deflection due to gravitation use the experimental

3) fact that a Solar system and all planar orbits are represented by a precessing conical section:

$$r = \frac{d}{1 + e \cos(x\theta)} \quad - (12)$$

for small  $x$ .

As in previous UFT papers such as UFT215 and UFT216 eq. (12) can be used in the equation:

$$\frac{d^2}{d\theta^2} \left( \frac{1}{r} \right) + \frac{1}{r} = - \frac{mr^2}{L} F(r) \quad - (13)$$

(Maria and Thornton eq. (7.21)) which is valid for all planar orbits in which total angular momentum is conserved. Eqs. (12) and (13)

give:

$$F(r) = - \frac{kx^2}{r} - \frac{k(1-x^2)d}{r^3} \quad - (14)$$

where  $k$  is a constant. If:

$$\underline{p} = p_r \underline{e}_r \quad - (15)$$

$$\underline{\omega} = \omega_r \underline{e}_r \quad - (16)$$

4) then:

$$F = \frac{dp_r}{dt} - \phi \omega_r = -\frac{kx^2}{r^2} - \frac{k(1-x^2)d}{r^3} \quad - (17)$$

For small deviation from a Newtonian orbit:

$$\frac{dp_r}{dt} = -\frac{kx^2}{r^2}, \quad - (18)$$

$$x \sim 1 \quad - (19)$$

to an excellent approximation. Eq. (18) is the equivalence principle. So:

$$\phi \omega_r = \frac{k(1-x^2)d}{r^3} \quad - (20)$$

to an excellent approximation, where:

$$\phi = -\frac{k}{r} \quad - (21)$$

and

$$k = mM\bar{G}. \quad - (22)$$

So:

$$\omega_r = -\left(1-x^2\right) \frac{d}{r^2} \quad - (23)$$

where

$$d = \frac{b^2}{a} \quad - (24)$$

5) Here  $d$  is the half right latitude,  $a$  and  $b$  are the semi major and minor axes, and  $e$  is the eccentricity. In light deflection by the sun the orbit is a hyperbola. The total deflection for the hyperbola is:

$$\Delta\phi = 2 \sin^{-1} \frac{1}{e} \quad - (25)$$

As shown in UFT 216, for small angles of deflection at closest approach,  $R_0$ :

$$\sin \phi \sim \phi = \frac{1}{e} = \left[ \frac{m^2 d R_0}{c^2 L^2} \left( v^2 - \frac{L^2}{m^2} \left( \frac{c^2 - 1}{R_0^2} \right) \right) - \frac{1}{e} \right]^{-1} \quad - (26)$$

where  $v$  is the velocity of a mass  $m$  orbiting a mass  $M$ . In the Newtonian limit this equation reduces to:

$$\sin \phi \sim \phi = \frac{1}{e} = \left( \frac{R_0 v^2}{MG} - 1 \right)^{-1} \quad - (27)$$

and for a photon of mass  $m$ :

$$v \rightarrow c \quad - (28)$$

$$\text{so} \quad \Delta\phi = \frac{2MG}{R_0 c^2} \quad - (29)$$

This is the so called Newtonian

result for light deflection due to gravitation.

By experimental observation the value of light deflection by any mass  $M$  is:

$$2\alpha = \frac{4MG}{R_0 c^2} \quad (30)$$

The experimental result is explained by a choice of  $\alpha$  in eq. (26) and therefore by a choice of spin connection  $\omega_r$  from eq. (23). If eq. (28) is assumed the correction needed to produce eq. (30) is:

$$\frac{m^2 d R_0 c^2}{L^2} = \frac{R_0 c^2}{mG}$$

$$\rightarrow \frac{R_0 c^2}{mG} - \frac{m^2 d R_0}{x^2 m^2} \left( \frac{x^2 - 1}{R_0^2} \right)$$

$$= \frac{R_0 c^2}{mG} - \frac{d}{R_0} \left( \frac{x^2 - 1}{x^2} \right)$$

$$= \frac{R_0 c^2}{mG} + \frac{d}{R_0} \left( \frac{1 - x^2}{x^2} \right) \quad (30)$$

7) Now use Q's result:

$$d = R_0(1 + \epsilon) - (31)$$

(Mauria and Thornton eq. (7.44)) to find

Ans: 
$$2\phi = \frac{2R_0 c^2}{M G} + 2(1 + \epsilon) \left( \frac{1 - x^2}{x^2} \right) - (31)$$

Experimentally:

$$(1 + \epsilon) \left( \frac{1 - x^2}{x^2} \right) = \frac{R_0 c^2}{M G} - (32)$$

The eccentricity  $\epsilon$  is given by Q's derived.  
angle of deflection is eq. (25):

$$\frac{1}{\epsilon} = \sin \left( \frac{\Delta\phi}{2} \right) - (33)$$

For small deflections:

$$\frac{1}{\epsilon} \sim \frac{\Delta\phi}{2} - (34)$$

$$\text{so } \left( 1 + \frac{2}{\Delta\phi} \right) \left( \frac{1 - x^2}{x^2} \right) = \frac{R_0 c^2}{M G} - (35)$$

For small deflections:

$$x \sim 1 - (36)$$

8) so

$$1 - x^2 = \frac{R_{oc}^2}{\underline{m}G} \left( 1 + \frac{2}{\Delta\psi} \right)^{-1} - (37)$$

However:  $\Delta\phi = \frac{4 R_{oc}^2}{mG} - (38)$

so  $1 - x^2 = \frac{\Delta\phi}{4} \left( 1 + \frac{2}{\Delta\psi} \right)^{-1} - (39)$

so using eq. (23):

$$\omega_r = - \frac{\Delta\phi}{4} \left( 1 + \frac{2}{\Delta\psi} \right)^{-1} \frac{d}{r^2} - (40)$$

Finally from eq. (31):

$$\begin{aligned} d &= R_o (1 + \epsilon) \\ &= R_o \left( 1 + \frac{2}{\Delta\psi} \right) \end{aligned} - (41)$$

From eqs. (40) and (41):

$$\boxed{\omega_r = - \frac{\Delta\phi}{4} \frac{R_o}{r^2}} - (42)$$

i.e.



at closest approach: ---

$$r = R_0 \quad - (43)$$

So:

$$\boxed{\omega_r = - \frac{\Delta\phi}{4R_0}} \quad - (44)$$

This is a universal result because it is found experimentally that light deflection due to gravitation is always given by twice the Newtonian value. Both  $\Delta\phi$  and  $R_0$  we know experimentally so  $\omega_r$  can be calculated and tabulated for any astronomical object.

It is named the  $\omega_r$  curvature for light deflection due to gravity.

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