

47(5): Orbital Precession is a Lorentz Force Equation
 Consider the minimal prescription in electrodynamics:

$$\underline{p} \rightarrow \underline{p} + e \underline{A} \quad - (1)$$

where \underline{p} is the momentum and \underline{A} the vector potential. Here e is the modulus of the charge of the electron. As in Note 47(2) the gravitomagnetic minimal prescription is:

$$\underline{p} \rightarrow \underline{p} + m \underline{v}_g \quad - (2)$$

where the gravitomagnetic vector potential is:

$$\underline{W}_g = \underline{v}_g \quad - (3)$$

The free particle Hamiltonian is:

$$H = \frac{1}{2m} (\underline{p} + m \underline{v}_g) \cdot (\underline{p} + m \underline{v}_g) \quad - (4)$$

$$= \frac{\underline{p}^2}{2m} + \frac{1}{2} m \underline{v}_g^2 + \frac{1}{2} \underline{L} \cdot \underline{\Omega}_g$$

where

$$\underline{L} = \underline{p} \times \underline{r} \quad - (5)$$

and

$$\underline{\Omega}_g = \nabla \times \underline{v}_g \quad - (6)$$

is the gravitomagnetic field. The orbital precession frequency is the Larmor precession frequency:

$$\underline{\Omega} = \frac{1}{2} \underline{\Omega}_g \quad - (7)$$

$\underline{\Omega}$ is any experimentally measured precession frequency.

In the presence of a gravitational potential:

$$U(r) = - \frac{m M G}{r} \quad - (8)$$

2) The Hamiltonian becomes:

$$H = \frac{1}{2m} (\mathbf{p} + m\mathbf{v}_g) \cdot (\mathbf{p} + m\mathbf{v}_g) + \bar{U}(r) \quad - (9)$$

In the absence of a gravitomagnetic field the Hamiltonian reduces to:

$$H = \frac{1}{2m} \mathbf{p} \cdot \mathbf{p} + U(r) \quad - (10)$$

It is well known that the Hamiltonian (10) produces the conic section:

$$r = \frac{d}{1 + \epsilon \cos \theta} \quad - (11)$$

in plane polar coordinates (r, θ) , where d is the half right semi-latus and ϵ the eccentricity.

Therefore the gravitomagnetic field (6) produces the orbital precession frequency Ω , with Hamiltonian:

$$H = \frac{\mathbf{p}^2}{2m} + \frac{1}{2} m \mathbf{v}_g^2 + \frac{1}{2} \mathbf{L} \cdot \frac{\Omega}{g} + \bar{U}(r) \quad - (12)$$

The Lagrangian is calculated from the Hamiltonian using the canonical momentum:

$$\mathbf{p}_q = \frac{\partial L}{\partial \dot{\mathbf{q}}} \quad - (13)$$

Denote:

$$\dot{\mathbf{r}} = \frac{1}{m} (\mathbf{p} + m\mathbf{v}_g) \quad - (14)$$

then:

$$H = \mathbf{p} \cdot \dot{\mathbf{r}} - L \quad - (15)$$

The Lagrangian is therefore:

$$3) \mathcal{L} = \frac{1}{2} m (\underline{p} + m \underline{v}_g) \cdot (\underline{p} + m \underline{v}_g) - U(r) - m \underline{\dot{r}} \cdot \underline{v}_g \quad (16)$$

and the relevant Euler Lagrange equation is:

$$\nabla \mathcal{L} = \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \underline{\dot{r}}} \right) \quad (17)$$

The left hand side of eq. (17) is:

$$\begin{aligned} \nabla \mathcal{L} &= \nabla \left(\frac{1}{2} m \underline{\dot{r}} \cdot \underline{\dot{r}} - U(r) \right) - m \nabla (\underline{\dot{r}} \cdot \underline{v}_g) \quad (18) \\ &= -\nabla U(r) - m \nabla (\underline{\dot{r}} \cdot \underline{v}_g) \end{aligned}$$

In general:

$$\begin{aligned} \nabla (\underline{\dot{r}} \cdot \underline{v}_g) &= (\underline{\dot{r}} \cdot \nabla) \underline{v}_g + (\underline{v}_g \cdot \nabla) \underline{\dot{r}} \quad (19) \\ &\quad + \underline{\dot{r}} \times (\nabla \times \underline{v}_g) + \underline{v}_g \times (\nabla \times \underline{\dot{r}}) \\ &= (\underline{\dot{r}} \cdot \nabla) \underline{v}_g + \underline{\dot{r}} \times (\nabla \times \underline{v}_g) \end{aligned}$$

if it is assumed that:

$$(\underline{v}_g \cdot \nabla) \underline{\dot{r}} = \underline{0} \quad (20)$$

$$\nabla \times \underline{\dot{r}} = \underline{0} \quad (21)$$

and

$$\text{So:} \quad \nabla \mathcal{L} = -\nabla U(r) - m \left((\underline{\dot{r}} \cdot \nabla) \underline{v}_g + \underline{\dot{r}} \times (\nabla \times \underline{v}_g) \right) \quad (22)$$

The right hand side of eq. (17) is:

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\underline{r}}} \right) = \frac{d}{dt} (m \underline{r} - m \underline{v}_g) \quad - (23)$$

$$= m \ddot{\underline{r}} - m \frac{d \underline{v}_g}{dt}$$

In order to evaluate $\frac{d}{dt}$ total derivative of \underline{v}_g , consider one component, so for example:

$$\frac{d v_{gx}}{dt} = \frac{\partial v_{gx}}{\partial t} + \left(\frac{dx}{dt} \right) \left(\frac{\partial v_{gx}}{\partial x} \right) + \dots$$

- (24)

to first order.

Therefore in three dimensions:

$$\frac{d \underline{v}_g}{dt} = \frac{\partial \underline{v}_g}{\partial t} + (\dot{\underline{r}} \cdot \underline{\nabla}) \underline{v}_g \quad - (25)$$

Therefore:

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\underline{r}}} \right) = m \ddot{\underline{r}} - m \left(\frac{\partial \underline{v}_g}{\partial t} + (\dot{\underline{r}} \cdot \underline{\nabla}) \underline{v}_g \right) \quad - (26)$$

The Euler Lagrange equation is therefore:

$$-\underline{\nabla} U(r) - m \dot{\underline{r}} \times (\underline{\nabla} \times \underline{v}_g) = m \ddot{\underline{r}} - m \frac{\partial \underline{v}_g}{\partial t} \quad - (27)$$

$$\text{i.e.} \quad m \ddot{\underline{r}} = -\underline{\nabla} U(r) + m \frac{\partial \underline{v}_g}{\partial t} - m \dot{\underline{r}} \times (\underline{\nabla} \times \underline{v}_g)$$

- (28)

5) Now write $m \phi_g = -U(r)$, - (29)

and

$$\underline{\Omega}_g = \underline{\dot{r}} \times \underline{v}_g \quad - (30)$$

to find the gravitomagnetic Lorentz force law

$$\underline{F} = m \underline{\ddot{r}} = -m \left(\underline{E}_g + \underline{\dot{r}} \times \underline{\Omega}_g \right) \quad - (31)$$

where

$$\underline{E}_g = -\underline{\nabla} \phi_g - \frac{\partial \underline{v}_g}{\partial t} \quad - (32)$$

The precession of an orbit by the Lorentz force law (31) with precession frequency $\underline{\Omega}$ is therefore governed by

$$\underline{\Omega} = \frac{1}{2} \underline{\Omega}_g \quad - (33)$$

The Hamiltonian corresponding to eq. (31) is:

$$H = \frac{1}{2m} (\underline{p} + m \underline{v}_g) \cdot (\underline{p} + m \underline{v}_g) + U(r) \quad - (34)$$

and the gravitational potential energy:

$$U(r) = -\frac{GMm}{r} \quad - (35)$$

so from eqs. (29) and (35):

$$\phi_g = -\frac{MG}{r} \quad - (36)$$

It follows that

$$\underline{\nabla} \phi_g = -\frac{MG}{r^2} \underline{e}_r \quad - (37)$$

b) so
$$E_g = \frac{MG}{r^2} \underline{e}_r - \frac{\partial \psi_g}{\partial t} \quad - (38)$$

and
$$\underline{F} = m \underline{\ddot{r}} - \frac{mMG}{r^2} \underline{e}_r + m \frac{\partial \psi_g}{\partial t} - m \underline{\dot{r}} \times \underline{\Omega_g} \quad - (39)$$

where
$$\underline{\dot{r}} = \frac{1}{m} (\underline{p} + m \underline{v}_g) \quad - (40)$$

In the absence of the gravitomagnetic field, eq. (39) reduces to:

$$\underline{F} = m \underline{\ddot{r}} - \frac{mMG}{r^2} \underline{e}_r \quad - (41)$$

In case equations for a planar orbit:

$$\underline{v} = \underline{\dot{r}} = \frac{dr}{dt} \underline{e}_r = \frac{d}{dt} (r \underline{e}_r) \quad - (42)$$

$$= \dot{r} \underline{e}_r + r \dot{\theta} \underline{e}_\theta$$

in the absence of a gravitomagnetic field. The acceleration in the absence of a gravitomagnetic field is:

$$\underline{a} = \frac{d\underline{v}}{dt} = \frac{d}{dt} (\dot{r} \underline{e}_r + r \dot{\theta} \underline{e}_\theta) = \underline{\ddot{r}} \quad - (43)$$

Eq. (43) gives rise to the well known Coriolis and centripetal forces, as it provides 4FT papers. So the gravitomagnetic forces occur in addition to these forces.

7) The gravitomagnetic field is governed by the law:

$$\nabla \times \underline{\Omega}_g = \frac{4\pi G}{c^2} \underline{J}_g \quad - (44)$$

analogous to the Ampère law:

$$\nabla \times \underline{B} = \mu_0 \underline{J} \quad - (45)$$

in electrodynamics and magnetostatics. The vacuum gravitomagnetic permeability is:

$$\mu_{g0} = \frac{4\pi G}{c^2} \quad - (46)$$

The gravitomagnetic four potential is:

$$\underline{W}^{\mu} = \left(\phi_g, c \underline{W}_g \right) \quad - (47)$$

where

$$\underline{W}_g = \underline{v}_g \quad - (48)$$