

SS(3): Hydrodynamics and Fluid Dynamics from ECE Wave Equation.

In previous work it has been shown that fluid dynamics is reduced to one wave equation:

$$\square v^\mu = \frac{1}{a_0^2} T^\mu \quad - (1)$$

Let

$$v^\mu = \left(\frac{\Phi}{a_0}, \underline{v} \right) \quad - (2)$$

and

$$T^\mu = (a_0 q, \underline{T}) \quad - (3)$$

Here a_0 is the speed of sound. Eq. (1) reduces to the ECE wave equation:

$$(\square + R) v^\mu = 0 \quad - (4)$$

provided that

$$R v^\mu = -\frac{1}{a_0^2} T^\mu \quad - (5)$$

i.e

$$R \Phi = -q \quad - (6)$$

and

$$R \underline{v} = -\frac{1}{a_0^2} \underline{T} \quad - (7)$$

The units are:

$$\Phi = m^2 s^{-2}; \quad q = s^{-2}; \quad R = m^{-2}; \quad \underline{T} = m s^{-3}$$

$$\text{and } \underline{v} = m s^{-1}.$$

From eqs. (6) and (7):

$$\frac{\underline{v}}{\Phi} = \frac{a_0^2 \underline{T}}{q} \quad - (8)$$

2) so it is possible to find $\underline{\underline{E}}$ from eq. (8) and the curvature from eqs. (6) and (7). The existence of curvature means that the theory is developed in a space with finite torsion and curvature.

The scheme of calculation is to start with the vorticity equation derived from Navier Stokes equation:

$$\frac{\partial \underline{\underline{w}}}{\partial t} = \underline{\underline{v}} \times (\underline{\underline{v}} \times \underline{\underline{w}}) + \frac{1}{\rho} \underline{\underline{v}} \times \underline{\underline{v}} \times \underline{\underline{P}} + \frac{\mu}{\rho} \nabla^2 \underline{\underline{w}} \quad (9)$$

where the baroclinic term is:

$$\frac{1}{\rho} \underline{\underline{v}} \times \underline{\underline{v}} \times \underline{\underline{P}} = - \underline{\underline{v}} \times \left(\frac{1}{\rho} \underline{\underline{v}} \times \underline{\underline{P}} \right) = - \underline{\underline{v}} \times \underline{\underline{v}} \times \underline{\underline{h}} \quad (10)$$

is 0 in Kramé theory. So:

$$\frac{\partial}{\partial t} (\underline{\underline{v}} \times \underline{\underline{v}}) = \underline{\underline{v}} \times (\underline{\underline{v}} \times \underline{\underline{w}}) + \frac{1}{R} \nabla^2 \underline{\underline{w}} \quad (11)$$

where the Reynolds number is approximately:

$$R = \rho / \mu \quad (12)$$

$$\text{So:} \quad \frac{\partial \underline{\underline{w}}}{\partial t} + \underline{\underline{v}} \times (\underline{\underline{w}} \times \underline{\underline{v}}) = \frac{1}{R} \nabla^2 \underline{\underline{w}} \quad (13)$$

Now use:

$$\begin{aligned} \nabla^2 \underline{\underline{w}} &= \underline{\underline{v}} (\underline{\underline{v}} \cdot \underline{\underline{w}}) - \underline{\underline{v}} \times (\underline{\underline{v}} \times \underline{\underline{w}}) \quad (14) \\ &= - \underline{\underline{v}} \times (\underline{\underline{v}} \times \underline{\underline{w}}) \end{aligned}$$

because: $\nabla \cdot \underline{w} = \nabla \cdot \nabla \times \underline{v} = 0 - (15)$

Therefore:

$$\frac{d}{dt} (\nabla \times \underline{v}) = \nabla \times (\underline{v} \times \underline{w}) - \frac{\mu}{\rho} \nabla \times (\nabla \times \underline{w}) - (16)$$

and $\frac{d\underline{v}}{dt} = \underline{v} \times \underline{w} - \frac{\mu}{\rho} \nabla \times \underline{w} - (17)$

as in previous work. So \underline{v} can be calculated from eq. (17). Note that:

$$\underline{v} \times (\nabla \times \underline{v}) = \frac{1}{2} \nabla v^2 - (\underline{v} \cdot \nabla) \underline{v} - (18)$$

and $\nabla \times (\nabla \times \underline{v}) = \nabla (\nabla \cdot \underline{v}) - \nabla^2 \underline{v} - (19)$

so:

$$\frac{D\underline{v}}{Dt} = \frac{d\underline{v}}{dt} + (\underline{v} \cdot \nabla) \underline{v} = \frac{1}{2} \nabla v^2 - \nabla (\nabla \cdot \underline{v}) + \nabla^2 \underline{v} - (20)$$

This is a useful form of the vorticity equation and resembles the Navier Stokes equation.

Eq. (1) is derived with the Lorenz condition:

$$\frac{1}{a_0^2} \frac{d\Phi}{dt} + \nabla \cdot \underline{v} = 0 - (21)$$

from which Φ can be calculated given \underline{v} from eq. (20). The charge and current are calculated

using:

$$\underline{q} = \underline{\nabla} \cdot \left(\left(\underline{v} \cdot \underline{\nabla} \right) \underline{v} \right) - (22)$$

and:

$$\underline{J} = a_0^2 \underline{\nabla} \times (\underline{\nabla} \times \underline{v}) - \frac{d}{dt} \left(\left(\underline{v} \cdot \underline{\nabla} \right) \underline{v} \right) - (23)$$

The $\underline{\Phi}$ calculated from eq. (21) must be self consistent wth eq. (8):

$$\underline{q} \cdot \underline{v} = a_0^2 \underline{\Phi} \cdot \underline{J} - (24)$$

So given \underline{v} , \underline{q} , \underline{J} and $\underline{\Phi}$, the speed of sound a_0 in a given system can be calculated.

Finally, the curvature of the system can be found from eqs. (6) & (7).

Note carefully that these equations are for the aether or vacuum and are worked out in terms of the potentials $\underline{\Phi}$ and \underline{v} . The vacuum transfers energy to a circuit.
