

57(3): Preparing Calculation for the Lamb Shift.

In this calculation we begin with the Dirac type Hamiltonian of UFT 253, eq. (41):

$$H\psi = \left( e\phi + mc^2 + \frac{1}{2m} \underline{\sigma} \cdot \underline{p} \underline{\sigma} \cdot \underline{p} + \frac{e}{4m^2 c^2} \underline{\sigma} \cdot \underline{p} \phi \underline{\sigma} \cdot \underline{p} \right) \psi \quad (1)$$

The relevant spin orbit term is:

$$H_1 \psi = \frac{e}{4m^2 c^2} (\underline{\sigma} \cdot \underline{p} \phi \underline{\sigma} \cdot \underline{p}) \psi \quad (2)$$

where

$$\underline{p} \psi = -i\hbar \underline{\nabla} \psi \quad (3)$$

so:

$$H_1 \psi = -\frac{ie\hbar}{4m^2 c^2} \underline{\sigma} \cdot \underline{\nabla} \phi \underline{\sigma} \cdot \underline{p} \psi \quad (4)$$

$$= -\frac{ie\hbar}{4m^2 c^2} \underline{\sigma} \cdot \underline{\nabla} (\phi \underline{\sigma} \cdot \underline{p} \psi)$$

$$= -\frac{ie\hbar}{4m^2 c^2} \left( \underline{\nabla} (\underline{\sigma} \cdot \underline{p}) \phi \psi + \underline{\sigma} \cdot \underline{p} \underline{\nabla} \cdot (\phi \psi) \right)$$

Now we:

$$\underline{\nabla} (\phi \psi) = \psi \underline{\nabla} \phi + \phi \underline{\nabla} \psi \quad (5)$$

so:

$$H_1 \psi = -\frac{ie\hbar}{4m^2 c^2} \left( \underline{\sigma} \cdot \underline{\nabla} \phi \underline{\sigma} \cdot \underline{p} \psi + \phi \underline{\sigma} \cdot \underline{\nabla} \psi \underline{\sigma} \cdot \underline{p} \right) \quad (6)$$

+ ...

2) The Coulomb potential is used as follows:

$$\phi = -\frac{e}{4\pi\epsilon_0 r} \quad (7)$$

So

$$\underline{\nabla} \phi = \frac{e}{4\pi\epsilon_0 r^3} \underline{r} \quad (8)$$

and

$$H_1 \psi = -\frac{ie^2 \hbar}{16\pi\epsilon_0 m^2 c^2} \left( \frac{1}{r^3} \underline{\sigma} \cdot \underline{r} \underline{\sigma} \cdot \underline{p} \psi - \frac{1}{r} \underline{\sigma} \cdot \underline{\nabla} \psi \underline{\sigma} \cdot \underline{p} \right) \quad (9)$$

Now use the Pauli algebra:

$$\begin{aligned} \underline{\sigma} \cdot \underline{r} \underline{\sigma} \cdot \underline{p} &= \underline{r} \cdot \underline{p} + i \underline{\sigma} \cdot \underline{r} \times \underline{p} \quad (10) \\ &= \underline{r} \cdot \underline{p} + i \underline{\sigma} \cdot \underline{L} \end{aligned}$$

where the orbital angular momentum is:

$$\underline{L} = \underline{r} \times \underline{p} \quad (11)$$

So:

$$Re(H_1 \psi) = \frac{e^2 \hbar}{16\pi\epsilon_0 m^2 c^2 r^3} \underline{\sigma} \cdot \underline{L} \psi + \dots \quad (12)$$

where

$$\underline{\sigma} \cdot \underline{L} = \int \psi^* \underline{\hat{\sigma}} \cdot \underline{\hat{L}} \psi d\tau \quad (13)$$

The spin angular momentum operator is:

$$\underline{\hat{S}} = \frac{\hbar}{2} \underline{\hat{\sigma}} \quad (14)$$

so

$$\underline{\hat{\sigma}} \cdot \underline{\hat{L}} = \frac{2}{\hbar} \underline{\hat{S}} \cdot \underline{\hat{L}} \quad (15)$$

In an representation:

$$\underline{\hat{S}} \cdot \underline{\hat{L}} \psi = \frac{\hbar^2}{2} (j(j+1) - l(l+1) - s(s+1)) \psi \quad (16)$$

where  $j$ ,  $l$  and  $s$  are quantum numbers.

Therefore:

$$\begin{aligned} \underline{\sigma} \cdot \underline{L} &= \langle \underline{\hat{\sigma}} \cdot \underline{\hat{L}} \rangle \\ &= \hbar (j(j+1) - l(l+1) - s(s+1)) \int \psi^* \psi d\tau \\ &= \hbar (j(j+1) - l(l+1) - s(s+1)) \quad (17) \end{aligned}$$

Therefore:

$$\text{Re}(\hat{H}_1 \psi) = \frac{e^2 \hbar^2}{16\pi \epsilon_0 m^2 c^2} (j(j+1) - l(l+1) - s(s+1)) \int \frac{\psi \psi^*}{r^3} d\tau \quad (18)$$

The Lamb shift is the effect of the vacuum or radiative or spin-orbit on this calculation. In the H atom:

$$4) \quad \int \frac{\psi \psi^*}{r^3} d\tau = \frac{1}{a_0^3 n^3 l(l+\frac{1}{2})(l+1)} \quad - (19)$$

where  $a_0 = \frac{4\pi\epsilon_0 \hbar^2}{me^4} \quad - (20)$

is the Bohr radius.

The Lamb shift is a difference of about a gigahertz between the  $2S_{1/2}$  and  $2P_{1/2}$  levels of the H atom. This is not predicted by the above Dirac theory. The  $2S_{1/2}$  level is about a gigahertz higher than the  $2P_{1/2}$  level.

The next note will begin to develop an explanation for this radiative correction using fluid dynamics and the potential of fluid spacetime.

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