

74(5) : Spin Coaxial Resonance and Magnetic Motors.

space-time when the magnetic motor is static Here is a balance described by :

$$\nabla \times \underline{A}^{(2)} = ig \underline{A}^{(2)} \times \underline{A}^{(3)} - (1)$$
$$= -\kappa_1 \underline{A}^{(2)}$$

where: $\underline{A}^{(2)} = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{-i\phi} - (2)$

$$\underline{A}^{(3)} = A^{(0)} \underline{k} - (3)$$

Here $A^{(0)}$ is produced by the magnetic field of the motor. The balance condition is therefore:

$$\nabla^2 \underline{A}^{(2)} + \kappa_1^2 \underline{A}^{(2)} = \underline{0} - (4)$$

which is a Helmholtz wave equation.

Similarly:

$$\nabla^2 \underline{A}^{(1)} + \kappa_1^2 \underline{A}^{(1)} = \underline{0} - (5)$$

and

$$\nabla^2 \underline{A}^{(3)} + \kappa_2^2 \underline{A}^{(3)} = \underline{0} - (6)$$

where $\kappa_2 = 0 - (7)$

2) The balance is broken when the right hand side of eqns. (4) to (6) are finite. If it is assumed that eq. (6), for example, becomes:

$$\frac{d^2 A_z}{dz^2} + \kappa_2^2 A_z = \mu_0 J_0 \cos(\kappa z) \quad \text{--- (8)}$$

then:

$$A_z = \frac{\mu_0 J_0 \cos(\kappa z)}{\kappa^2 - \kappa_2^2} \quad \text{--- (9)}$$

and when:

$$\kappa = \kappa_2 \quad \text{--- (10)}$$

then

$$A_z \rightarrow \infty \quad \text{--- (11)}$$

The torque produced in the magnetic assembly is:

$$\underline{T}_q = -A_z \underline{m} \times \underline{k} \quad \text{--- (12)}$$

and becomes very large. The property $J_0 \cos(\kappa z)$ is a property of the magnetic assembly and $\kappa_2 = 1/r_2$

is a spiral constant of spacetime. Under the resonance condition (10) the magnet...