

6(5) Basic Definition of Angular Momentum as a Tetrad.

The angular momentum was defined in paper JSS as a tetrad:

$$J^a_{\mu} = J^{(a)}_{\nu} g^{\nu\mu} \quad - (1)$$

which implies that there exist vectors W^a and V^{μ} :

$$W^a = J^a_{\mu} V^{\mu} \quad - (2)$$

i.e.

$$\begin{bmatrix} W^1 \\ W^2 \\ W^3 \end{bmatrix} = \begin{bmatrix} J^1_1 & J^1_2 & J^1_3 \\ J^2_1 & J^2_2 & J^2_3 \\ J^3_1 & J^3_2 & J^3_3 \end{bmatrix} \begin{bmatrix} V^1 \\ V^2 \\ V^3 \end{bmatrix} \quad - (3)$$

if we restrict consideration to three space dimensions.

Thus:

$$\begin{aligned} W^1 &= J^1_1 V^1 + J^1_2 V^2 + J^1_3 V^3 \\ W^2 &= J^2_1 V^1 + J^2_2 V^2 + J^2_3 V^3 \\ W^3 &= J^3_1 V^1 + J^3_2 V^2 + J^3_3 V^3 \end{aligned} \quad - (4)$$

It is known that angular momentum must be an antisymmetric tensor so:

$$\left. \begin{aligned} W^1 &= J^1_2 V^2 + J^1_3 V^3 \\ W^2 &= J^2_1 V^1 + J^2_3 V^3 \\ W^3 &= J^3_1 V^1 + J^3_2 V^2 \end{aligned} \right\} \quad - (5)$$

In 3-D space we can write this as:

$$\left. \begin{aligned} W_1 &= J_{12} V_2 + J_{13} V_3 \\ W_2 &= J_{21} V_1 + J_{23} V_3 \\ W_3 &= J_{31} V_1 + J_{32} V_2 \end{aligned} \right\} - (6)$$

where: $J_{ij} = \frac{1}{2} \epsilon_{ijk} J_k$ - (7)

$$\left. \begin{aligned} W_1 &= J_3 V_2 - J_2 V_3 \\ W_2 &= J_3 V_1 - J_1 V_3 \\ W_3 &= J_1 V_2 - J_2 V_1 \end{aligned} \right\} - (8)$$

It is seen that:

$$\begin{aligned} \underline{W} &= \underline{J} \times \underline{V} - (9) \\ &= \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ J_1 & J_2 & J_3 \\ V_1 & V_2 & V_3 \end{vmatrix} \end{aligned}$$

Now identify:

$$\underline{W} = \underline{P}, \quad \underline{V} = \underline{r} / r^2 - (10)$$

$$\underline{P} = \underline{J} \times \underline{r} / r^2 - (11)$$

i.e. $\boxed{\underline{J} = \underline{r} \times \underline{P}}$ - (12)

3)

From paper 55, :

$$\underline{N}^a(\text{torque}) = c \underline{J}^{(0)} \underline{T}^a(\text{torsion}) - (13)$$

and

$$\underline{N}^a = c (d \wedge \underline{J}^a + \omega^a{}_b \wedge \underline{J}^b) - (14)$$

This is the definition of torque (in joules) in a spacetime which is general but has no zero torsion and curvature.

In paper 76, a constant torsion model is used :

$$\underline{N}^a = \text{constant} - (15)$$

$$d \wedge \underline{N}^a = 0 - (16)$$

with $\omega^a{}_b = -\frac{1}{2} \kappa \epsilon^a{}_{bc} \underline{v}^c - (17)$

Adopting the complex circular basis :

$$\underline{N}^{(1)*} = c \left(\underline{\nabla} \times \underline{J}^{(1)*} - i \kappa \underline{v}^{(2)} \times \underline{J}^{(3)} \right) - (18)$$

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and the constant torque is :

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$$\begin{aligned}
 \underline{N}^{(3)*} &= -i \kappa c q \underline{v}^{(1)} \times \underline{J}^{(2)} \\
 &= c \kappa \underline{J}^{(0)} \underline{k} \\
 &= c \kappa m r v \underline{k}
 \end{aligned}
 \quad \left. \vphantom{\underline{N}^{(3)*}} \right\} \text{--- (19)}$$

Thus:

$$\boxed{\underline{N} = c m v \underline{k}} \quad \text{--- (20)}$$

because:

$$\kappa r = 1. \quad \text{--- (21)}$$

This is a precise analogy of $\underline{B}^{(3)}$ theory
in electrodynamics.

As shown in notes 76 (5),

$$\left. \begin{aligned}
 \theta &= v_0 \tau / r \\
 r &= v_0 \tau / \theta
 \end{aligned} \right\} \text{--- (22)}$$

from the angular momentum:

$$\omega = \frac{d\theta}{dt} = \frac{v}{r} \quad \text{--- (23)}$$

with $v = v_0 = \text{constant}$. Eq (22) is the
 equation of the hyperbolic spiral.