

79(4) : Lorentz Force Equations from ECE Theory.

This is derived from the transformation properties of the Carter tensor. As in eq. (25) of paper 71, the general mixed index tensor transforms as:

$$T^{a'\mu'}_{b'\nu'} = \Lambda^{a'}_a \frac{dx^\mu}{dx^{a'}} \Lambda^{b'}_b \frac{dx^\nu}{dx^{b'}} T^{ab}_{\mu\nu} \quad - (1)$$

so the Carter tensor transforms as:

$$T^{a'}_{\mu'} = \Lambda^{a'}_a \frac{dx^\mu}{dx^{a'}} \frac{dx^\nu}{dx^{b'}} T^a_{\mu\nu} \quad - (2)$$

Here $\Lambda^{a'}_a$ denotes the Lorentz transformation and $dx^\mu/dx^{a'}$ and $dx^\nu/dx^{b'}$ denote the general coordinate transformation, which is the limit of Minkowski spacetime become Lorentz transformations.

In the special case where rotation and translation are mutually independent, the ECE field equations are:

$$d \wedge F^a_{\mu\nu} = 0 \quad - (3)$$

$$d \wedge \tilde{F}^a_{\mu\nu} = \mu_0 J^a_{\mu\nu} \quad - (4)$$

For each polarization index a these are the Maxwell Heaviside (MH) field equations:

$$d \wedge F_{\mu\nu} = 0 \quad - (5)$$

$$d \wedge \tilde{F}_{\mu\nu} = \mu_0 J_{\mu\nu} \quad - (6)$$

It is well known that $F_{\mu\nu}$ in eqn. (5) transforms as:

$$F_{\mu'\nu'} = \Lambda^{\mu}_{\mu'} \Lambda^{\nu}_{\nu'} F_{\mu\nu} \quad - (7)$$

for each state of polarization a . Eq. (7) is the limit of eq. (2) where:

$$\frac{dx^{\mu}}{dx^{\mu'}} \rightarrow \Lambda^{\mu}_{\mu'} ; \quad \frac{dx^{\nu}}{dx^{\nu'}} \rightarrow \Lambda^{\nu}_{\nu'} \quad - (8)$$

for each fixed a .

The Lorentz free equation is obtained from eq. (7) as follows. From eq. (7):

$$\begin{aligned} \underline{E}' &= \gamma (\underline{E} + \underline{v} \times \underline{B}) - \dots \text{higher order terms} \\ \underline{B}' &= \gamma (\underline{B} - \frac{\underline{v}}{c^2} \times \underline{E}) - \dots \text{"} \end{aligned} \quad - (9)$$

and

$$\underline{F}' = e \gamma (\underline{E} + \underline{v} \times \underline{B}) - \dots \quad - (10)$$

In the non-relativistic limit:

$$\gamma \rightarrow 1 \quad - (11)$$

$$\boxed{\underline{F}' = e (\underline{E} + \underline{v} \times \underline{B})} \quad - (12)$$

3) Hence for each sense of polarization:
 $a = (1), (2), (3) \quad \text{--- (13)}$

Here is a corresponding Lorentz force:

$$\underline{F}^{(1)'} = e \left(\underline{E}^{(1)} + \underline{v} \times \underline{B}^{(1)} \right) \quad \text{--- (14)}$$

$$\underline{F}^{(2)'} = e \left(\underline{E}^{(2)} + \underline{v} \times \underline{B}^{(2)} \right) \quad \text{--- (15)}$$

$$\underline{F}^{(3)'} = e \left(\underline{E}^{(3)} + \underline{v} \times \underline{B}^{(3)} \right) \quad \text{--- (16)}$$

\underline{I}_L eq. (16) there is no radiated $\underline{E}^{(3)}$, so:

$$\underline{F}^{(3)'} = e \underline{v} \times \underline{B}^{(3)}$$

$$= -ie g \underline{v} \times (\underline{A}^{(1)} \times \underline{A}^{(2)}) \quad \text{--- (17)}$$

For a plane wave:

$$\underline{A}^{(1)} = \underline{A}^{(2)*} = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} - i\underline{j}) e^{i\phi} \quad \text{--- (18)}$$

For a static magnetic field of \underline{B} standard model:

$$\underline{F} = e \underline{v} \times (\underline{v} \times \underline{A}) \quad \text{--- (19)}$$

Eqs. (14) and (15) are first order forces and eq. (16) is a second order force.

+) These first and second order forms are also present
 in the Hamilton-Jacobi method.

Due to the Lorentz invariance of the B

Cyclic theorem:

$$\underline{B}^{(1)} \times \underline{B}^{(2)} = i \underline{B}^{(3)*} \quad \text{--- (20)}$$

et cyclicum
 The polarizations (1), (2) and (3) are
 invariant under a frame transformation. The
 Lorentz invariance of eq. (20) is obvious from
 the fact that it is:

$$\underline{e}^{(1)} \times \underline{e}^{(2)} = i \underline{e}^{(3)*} \quad \text{--- (21)}$$

where:

$$\underline{e}^{(1)} = \underline{e}^{(2)*} = \frac{1}{\sqrt{2}} (\underline{i} - i \underline{j})$$

$$\underline{e}^{(3)} = \underline{e}^{(3)*} = \underline{k} \quad \text{--- (22)}$$

The unit vectors \underline{i} , \underline{j} and \underline{k} are not affected
 by a Lorentz transformation. For example

relativistic momentum is:

$$\underline{p} = \gamma m \underline{v} \quad \text{--- (23)}$$

e.g.

$$p_z \underline{k} = \gamma m v_z \underline{k} \quad \text{--- (24)}$$

and \underline{k} is the same.