

90(3) now renumbered 92(3) : Properties of Spi Comets

In note 90(1) it was shown that the basic equation

$$\frac{d^2 \phi}{dr^2} + \left(\frac{2}{r} + \omega_r \right) \frac{d\phi}{dr} + \frac{\phi}{r^2} \left(2r\omega_r + r^2 \frac{d\omega_r}{dr} \right) = -\frac{f}{\epsilon_0} \quad \text{--- (1)}$$

reduces to:

$$\frac{d^2 \phi}{dr^2} + 2\beta \frac{d\phi}{dr} + \kappa_0^2 \phi = -\frac{f}{\epsilon_0} \quad \text{--- (2)}$$

where:

$$\omega_r = \kappa_0^2 r - 4\beta \log_e r - \frac{4}{r} \quad \text{--- (3)}$$

a) \underline{I}_h limit:

$$\kappa_0 \rightarrow 0, \beta \rightarrow 0 \quad \text{--- (4)}$$

it is seen that the wave dependence is recovered.

b) \underline{I}_c limit:

$$\beta \rightarrow 0, \kappa_0 \neq 0 \quad \text{--- (5)}$$

eqn. (2) becomes an undamped resonator equation:

$$\frac{d^2 \phi}{dr^2} + \kappa_0^2 \phi = -\frac{f}{\epsilon_0} \quad \text{--- (6)}$$

and $\phi \rightarrow \infty$ at resonance. \underline{I}_c limit

(5) $\omega_r \rightarrow \kappa_0^2 r - \frac{4}{r} \quad \text{--- (7)}$

2) From eq. (7) it is seen that there is a characteristic frequency ω_0 , which can be expressed as:

$$\kappa_0^2 = \frac{1}{r} \left(\omega_r + \frac{4}{r} \right) \quad - (8)$$

This is a Hooke's law ~~frequency~~ ^{wavenumber} defined by the spi connection, and the radial coordinate. For an undamped oscillator with a simple cosinusoidal dependence

$$\rho = -\rho_0 \cos(\kappa r) \quad - (9)$$

resonance occurs at:

$$\kappa_0 = \kappa \quad - (10)$$

CONCLUSION

The important deduction is that the potential ϕ can be amplified to infinity by eq. (10). The spi connection in this case

is:

$$\omega_r = \kappa^2 r - \frac{4}{r} \quad - (11)$$