

78(4) Rotation as a General Coordinate Transformation

The general coordinate transformation in general relativity generates the vector \tilde{V}_α from the vector V_μ as follows:

$$\tilde{V}_\alpha = \left(\frac{dx^\mu}{dx^{\tilde{\alpha}}} \right) V_\mu \quad - (1)$$

$$:= \epsilon^{\mu}_{\tilde{\alpha}} V_\mu. \quad - (2)$$

The Lorentz transformation is a special case of $\epsilon^{\mu}_{\tilde{\alpha}}$ and rotation is a special case of the Lorentz transformation. This is because the rotation group is a sub-group of the Lorentz group and the latter is a sub-group of the Einstein group. Therefore a general:

$$\epsilon^{\mu}_{\tilde{\alpha}} = \frac{dx^\mu}{dx^{\tilde{\alpha}}}. \quad - (3)$$

This generalizes the $\epsilon^{\mu}_{\tilde{\alpha}}$ used in Ryder's Chapter 6 to a spacetime which is general has torsion and curvature. Rotation is defined in this spacetime by:

$$\epsilon^{\rho\sigma} = -\epsilon^{\sigma\rho} = g^{\rho\kappa} \epsilon^{\sigma}_{\kappa} \quad - (4)$$

where the Kronecker delta is defined by:

$$\delta^{\mu}_{\sigma} = g^{\mu\nu} g_{\nu\sigma}. \quad - (5)$$

It is to be noted that the general coordinate transform (1) also applies to the partial derivatives

$$2) \quad \partial_\alpha = \left(\frac{\partial x^\mu}{\partial x^\alpha} \right) \partial_\mu \quad - (6)$$

and that the set of partial derivatives $\{\partial_\mu\}$ at a point p in the Lorentz manifold defines a basis set (Carroll, chapter 23).

In Ryder chapter 13 it is stated that

$$\text{if: } x^\mu = \epsilon^\mu_\nu x^\nu = X^\mu_{p^\sigma} \epsilon^{\sigma\rho} \quad - (7)$$

$$\text{then: } X^\mu_{p^\sigma} = \frac{1}{2} (\delta^\mu_\sigma - g^\mu_\sigma) \quad - (8)$$

is a Minkowski spacetime. This is the point at which the rank three tensor $X^\mu_{p^\sigma}$ is introduced into the theory of rotations, and leads to the definition of the rank three angular momentum density tensor. In Cartan's differential geometry, the Cartan tensor is also a rank three tensor, so its origin is rotational. The gamma connection also has three indices but is not a tensor. The Cartan tensor is the tensor:

$$T^\mu_{p^\sigma} = \Gamma^\mu_{p^\sigma} - \Gamma^\mu_{\sigma p} \quad - (9)$$

and the difference of two general gamma connections is a tensor (see Carroll for more details). Hence a proof of eqs. (7) and (8) is important. Since Ryder does not give a proof (the following) is suggested

3) First generalize Ryder's equation (7) to any vector in a space-time with torsion and curvature:

$$\nabla^\mu = \nabla_{\rho\sigma}^\mu \epsilon^{\rho\sigma} \quad - (10)$$

Where:

$$\epsilon^{\rho\sigma} = -\epsilon^{\sigma\rho} \quad - (11)$$

$$\nabla_{\rho\sigma}^\mu = -\nabla_{\sigma\rho}^\mu \quad - (12)$$

This means that any vector ∇^μ can be expressed as a rank three tensor of type (12). In 3-D

Euclidean space this result is well known, any axial vector can be expressed as a rank two antisymmetric tensor:

$$\nabla_i = \frac{1}{2} \epsilon_{ijk} \nabla_{jk} \quad - (13)$$

where ϵ_{ijk} is the rank three totally antisymmetric unit tensor in 3-D Euclidean space.

In general, eq. (10) is:

$$\frac{1}{2} \nabla^\mu = \nabla_{01}^\mu \epsilon^{01} + \nabla_{02}^\mu \epsilon^{02} + \nabla_{03}^\mu \epsilon^{03} + \nabla_{12}^\mu \epsilon^{12} + \nabla_{13}^\mu \epsilon^{13} + \nabla_{23}^\mu \epsilon^{23} \quad - (14)$$

This can be expressed as a 4-D matrix with a structure similar to the electromagnetic field

matrix:

4)

$$\bar{V}^{\mu}_{\rho\sigma} = \begin{bmatrix} 0 & \bar{V}^{\mu}_{01} & \bar{V}^{\mu}_{02} & \bar{V}^{\mu}_{03} \\ \bar{V}^{\mu}_{10} & 0 & \bar{V}^{\mu}_{12} & \bar{V}^{\mu}_{13} \\ \bar{V}^{\mu}_{20} & \bar{V}^{\mu}_{21} & 0 & \bar{V}^{\mu}_{23} \\ \bar{V}^{\mu}_{30} & \bar{V}^{\mu}_{31} & \bar{V}^{\mu}_{32} & 0 \end{bmatrix} \quad - (15)$$

Rotation About the 3 Axis ($Z=3$), Axial Vectors

This is the special case:

$$\mu = 3, \quad \epsilon^{12} = -\epsilon^{21} = 1 \quad - (16)$$

all other $\epsilon^{\mu\nu} = 0$, so:

$$\bar{V}^3 = 2\bar{V}^3_{12} \quad - (17)$$

This is the magnetic field relation used in paper 93 apart from factor 2. For a fixed 3 it is the same as eq (13).

Polar Vector Relations

These are:

$$\frac{1}{2} \bar{V}^i = \bar{V}^i_{0i}, \quad i = 1, 2, 3 \quad - (18)$$

Apart from a factor 2 they are the electric field relations used in paper 93.

We now define the rank three tensor:

$$\bar{V}^{\mu}_{\rho\sigma} = \bar{V}^{\mu} \epsilon_{\rho\sigma} \quad - (19)$$

1 where it general:

$$E_{\rho\sigma} = g_{\rho\kappa} E^{\kappa}_{\sigma} \quad - (20)$$

$$E^{\kappa}_{\sigma} = \frac{\partial x^{\kappa}}{\partial x^{\sigma}} \quad - (21)$$

The matrix E^{κ}_{σ} is similar to the definition we use for tetrad differential geometry:

$$\nabla^a = e^a_{\mu} \nabla^{\mu} \quad - (22)$$

So:

$$e^{\kappa}_{\sigma} = \frac{\partial x^{\kappa}}{\partial x^{\sigma}} \quad - (23)$$

is a special case where:

$$E^{\mu}_{\nu} = e^{\mu}_{\nu} = e^a_{\nu} e^{\mu}_a \quad - (24)$$

The Cartan-Kasner form is:

$$T^a_{\mu\nu} = d \wedge e^a_{\nu} + \omega^a_{\mu b} \wedge e^b_{\nu} \quad - (25)$$

$$= dx^{\mu} e^a_{\nu} - dx^{\nu} e^a_{\mu} + \omega^a_{\mu b} e^b_{\nu} - \omega^a_{\nu b} e^b_{\mu} \quad - (26)$$

and the Cartan-Kasner tensor is

$$T^{\kappa}_{\mu\nu} = \Gamma^{\kappa}_{\mu\nu} - \Gamma^{\kappa}_{\nu\mu} \quad - (27)$$

b) They are related by:

$$T_{\mu\nu}^a = \eta^a{}^\kappa T_{\mu\nu}^\kappa \quad - (28)$$

and the tetrad postulate:

$$D_\mu \eta^a{}_\nu = 0. \quad - (29)$$

(19) Comparison of the structures of eqs. (28) and (19) shows that the central tensor is:

$$T_{\rho\sigma}^a = \delta_\rho^\sigma \epsilon^a - \delta_\sigma^\rho \epsilon^a + \omega_{\rho b}^a \epsilon_\sigma^b - \omega_{\sigma b}^a \epsilon_\rho^b \quad - (30)$$

and derives from a structure:

$$T_{\rho\sigma}^\mu = \delta_{\rho\sigma}^\mu + \dots \quad - (31)$$

that reflects the fundamental geometry of 4-D spacetime, another expression of which is eq. (19).