

CORRECTIONS TO THE RAYLEIGH JEANS DENSITY OF STATES

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ABSTRACT

From elementary considerations the Rayleigh Jeans density of states is corrected with two terms left out by Rayleigh in 1900. These are higher order infinitesimal terms that contribute significantly to the derivation of the Stefan Boltzmann law from the Planck distribution. The dogmatic approach of the standard model is therefore incorrect by a large factor. The B field of ECE theory also contributes significantly in circularly polarized radiation.

Keywords: ECE theory, Rayleigh Jeans density of states, corrections to the derivation of the Stefan Boltzmann law from the Planck distribution.

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1. INTRODUCTION

Recently in this series of two hundred and ninety one papers to date {1 - 10} the Evans Morris effects have been analyzed with various theories of refraction and reflection. It has been shown that the dogmatic approach to the theory of reflection and refraction is trivially incorrect. Elementary considerations of geometry show that the incident, reflected and refracted frequencies cannot be the same in general. The Evans Morris effects are colour changes in laser beam reflected and refracted from various materials, as described on the diary or blog of www.aias.us. During the course of extending the analysis to polychromatic radiation it became necessary to correct the Rayleigh Jeans density of states of 1900 and 1905 for two terms left out by Rayleigh and by Jeans. These higher order infinitesimal terms are described in Section 2, where it is shown that they result in a large correction to the Rayleigh Jeans density of states. It is shown that the \underline{B} field of ECE theory also results in a large correction when the radiation is circularly polarized.

2. CORRECTIONS OF THE RAYLEIGH JEANS LAW AND PLANCK DISTRIBUTION

Consider the number of oscillators N in a volume V of radiation (the number density) in the original Rayleigh calculation of 1900:

$$\frac{N}{V} = \frac{\omega^3}{6c^3\pi^2} \quad \text{--- (1)}$$

where ω is the angular frequency of monochromatic radiation and where c is the vacuum speed of light. The infinitesimal of number density was defined by Rayleigh as:

$$\frac{dN}{V} = \frac{1}{6c^3\pi^2} \left((\omega+d\omega)^3 - \omega^3 \right) = \frac{\omega^2 d\omega}{2c^3\pi^2} + \frac{\omega (d\omega)^2}{2c^3\pi^2} + \frac{(d\omega)^3}{6c^3\pi^2} \quad \text{--- (2)}$$

and he doubled this result by assuming that the radiation contains two states of polarization.

Presumably these are right and left circularly polarized. Therefore:

$$\frac{dN}{V} = \frac{\omega^2}{\pi^2 c^3} d\omega + \frac{\omega}{\pi^2 c^3} (d\omega)^2 + \frac{(d\omega)^3}{3\pi^2 c^3} \quad - (3)$$

Rayleigh left out the second two terms of Eq. (3) without attempting to justify the approximation, so the dogmatic Rayleigh Jeans density of states is:

$$\frac{dN}{V} = \frac{\omega^2}{\pi^2 c^3} d\omega. \quad - (4)$$

The number of oscillators N in a radiation volume V of polychromatic radiation is found by integrating over frequencies:

$$\frac{N}{V} = \frac{1}{V} \int dN = \int \frac{\omega^2}{\pi^2 c^3} d\omega = \frac{\omega^3}{3\pi^2 c^3}. \quad - (5)$$

This is the dogmatic Rayleigh Jeans density of states. It is well known to be incorrect at high frequencies because it goes to infinity. It is also well known that it was corrected by the use of the 1899 quantum theory of Planck, in which the average energy of a radiation oscillator is:

$$\langle E \rangle = \left(\frac{x}{1-x} \right) h\omega \quad - (6)$$

where:

$$x = \exp\left(-\frac{h\omega}{kT}\right) \quad - (7)$$

under the condition:

$$h\omega < kT. \quad - (8)$$

The infinitesimal of energy density from this theory is:

$$\frac{dU}{V} = \langle E \rangle \frac{dN}{V} \quad - (9)$$

and the energy density of radiation in joules per cubic metre is found by integrating over

frequency:

$$\frac{\bar{U}}{V} = \frac{1}{V} \int dU = \int \langle E \rangle \frac{dN}{V} d\omega. \quad - (10)$$

Using the dogmatic assertion (4) of Rayleigh, it is found that:

$$\frac{\bar{U}}{V} = \int \frac{h\omega^3}{\pi^2 c^3} \left(\frac{x}{1-x} \right) d\omega. \quad - (11)$$

In order to derive the experimental Stefan Boltzmann law a black body was defined as

having frequencies in the range from zero to infinity, so:

$$\frac{\bar{U}}{V} = \int_0^{\infty} \frac{h\omega^3}{\pi^2 c^3} \left(\frac{x}{1-x} \right) d\omega = \left(\frac{\pi^2 h^4}{15 c^3 h^3} \right) T^4 \quad - (12)$$

This is the old Stefan Boltzmann law derived from the Planck distribution. The energy U per

volume V of radiation from a black body radiator is proportional to the fourth power of

temperature. The number of Planck / Einstein photons of energy $h\omega$ per volume V is,

from Eq. (11):

$$\frac{N_p}{V} = \int \frac{\langle E \rangle}{h\omega} \frac{dN}{V} d\omega = \int \frac{\omega^2}{\pi^2 c^3} \left(\frac{x}{1-x} \right) d\omega \quad - (13)$$

For a black body radiator:

$$\frac{N_p}{V} = \int_0^{\infty} \frac{\omega^2}{\pi^2 c^3} \left(\frac{x}{1-x} \right) d\omega = \left(\frac{2 \zeta(3)}{\pi^2} \left(\frac{h}{c h} \right)^3 \right) T^3 \quad - (14)$$

where

$$\zeta(3) = 1.20206 \quad - (15)$$

is the third zeta function defined by:

$$\int_0^{\infty} \left(\frac{x^{2n}}{e^x - 1} \right) dx = (2n)! \zeta(2n+1) \quad - (16)$$

In the old theory the number of photons N_p per volume V of a black body radiator is proportional to the cube of temperature.

Both calculations used the dogmatic Rayleigh Jeans density of states (4) uncritically, and these calculations have always been claimed to verify the old quantum theory. It is straightforward to show as follows that these claims are not true. The correct density of states is Eq. (3), in which the second two terms contribute more oscillators per unit volume as follows.

Write Eq. (3) as:

$$\frac{dN}{V} = \frac{1}{V} (dN_1 + dN_2 + dN_3) \quad - (17)$$

where:

$$\frac{dN_2}{V} = \frac{\omega}{\pi^2 c^3} (d\omega)^2 \quad - (18)$$

and:

$$\frac{dN_3}{V} = \frac{(d\omega)^3}{3\pi^2 c^3} \quad - (19)$$

with the original Rayleigh term as:

$$\frac{dN_1}{V} = \frac{\omega^2}{\pi^2 c^3} d\omega \quad - (20)$$

From Eq. (18):

$$\left(\frac{dN_2}{V}\right)^{1/2} = \left(\frac{\omega}{\pi^2 c^3}\right)^{1/2} d\omega. \quad (21)$$

From Eq. (19):

$$\left(\frac{dN_3}{V}\right)^{1/3} = \frac{d\omega}{(3\pi^2 c^3)^{1/3}}. \quad (22)$$

Integrating Eq. (21) over frequency:

$$\frac{1}{V^{1/2}} \int (dN_2)^{1/2} = \int \left(\frac{\omega}{\pi^2 c^3}\right)^{1/2} d\omega = \frac{2\omega^{3/2}}{3(\pi^2 c^3)^{1/2}} \quad (23)$$

so there exists a number density of oscillators:

$$\frac{N_A}{V} = \frac{1}{V} \left(\int (dN_2)^{1/2} \right)^2 = \frac{4\omega^3}{9\pi^2 c^3}. \quad (24)$$

Similarly, integrating Eq. (22):

$$\frac{1}{V^{1/3}} \int (dN_3)^{1/3} = \int \frac{d\omega}{(3\pi^2 c^3)^{1/3}} = \frac{\omega}{(3\pi^2 c^3)^{1/3}} \quad (25)$$

so there exists a second number density of oscillators:

$$\frac{N_B}{V} = \frac{1}{V} \left(\int (dN_3)^{1/3} \right)^3 = \frac{\omega^3}{3\pi^2 c^3}. \quad (26)$$

The total number of oscillators in a volume V of radiation is therefore:

$$\frac{N}{V} = \frac{1}{V} (N_1 + N_A + N_B) = \frac{10}{9} \frac{\omega^3}{\pi^2 c^3}. \quad (27)$$

The original Rayleigh result is:

$$\frac{N}{V} = \frac{\omega^3}{3\pi^2 c^3}. \quad (28)$$

So the correction factor is 10/3, more than a factor three. The correct result (27) must be

obtained by integration of:

$$\frac{dN}{V} = \frac{10}{3} \frac{\omega^2}{\pi^2 c^3} d\omega \quad (29)$$

so the corrected Rayleigh Jeans density of states is:

$$\frac{dN}{V} = \frac{\omega^2 d\omega}{2c^3 \pi^2} + \frac{\omega (d\omega)^2}{2c^3 \pi^2} + \frac{(d\omega)^3}{6c^3 \pi^2} = \frac{10}{3} \frac{\omega^2}{\pi^2 c^3} d\omega \quad (30)$$

This must be used in the calculation of the Stefan Boltzmann law as described already. The proportionality constant of the Stefan Boltzmann law is increased by 10/3.

In Note 291(2) the evaluation of the definite integral:

$$\bar{I} = \int_a^b (dx)^{1/2} \quad (31)$$

is considered using the Riemann integration technique, and it shown that:

$$\int (dx)^{1/2} = x^{1/2} \quad (32)$$

so

$$N_2 = N_A, \quad N_3 = N_B \quad (33)$$

Finally in Note 291(3) it is shown that the $\underline{B}^{(3)}$ field of ECE theory contributes a classical energy density

$$\frac{E_n}{V} = \frac{1}{\mu_0} \underline{B}^{(3)} \cdot \underline{B}^{(3)*} = \epsilon_0 E^{(0)2} \quad (34)$$

to circularly polarized radiation where μ_0 is the vacuum permeability, E is the amplitude of the electric field strength and ϵ_0 is the vacuum permittivity. This increases the energy density by a factor of 5/4, so in circularly polarized radiation the complete correction to the Rayleigh Jeans law is

$$f_c = \frac{5}{4} \cdot \frac{10}{9} = \frac{50}{36} = \frac{25}{18} - (35)$$

In amorphous black body radiation however the B(3) field cancels to zero because it changes sign between right and left handed circularly polarized components.

These corrections to the original Rayleigh Jeans density of states (4) mean that the experimental testing of the law must be rethought, because Rayleigh's derivation contains unjustifiable approximations.

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